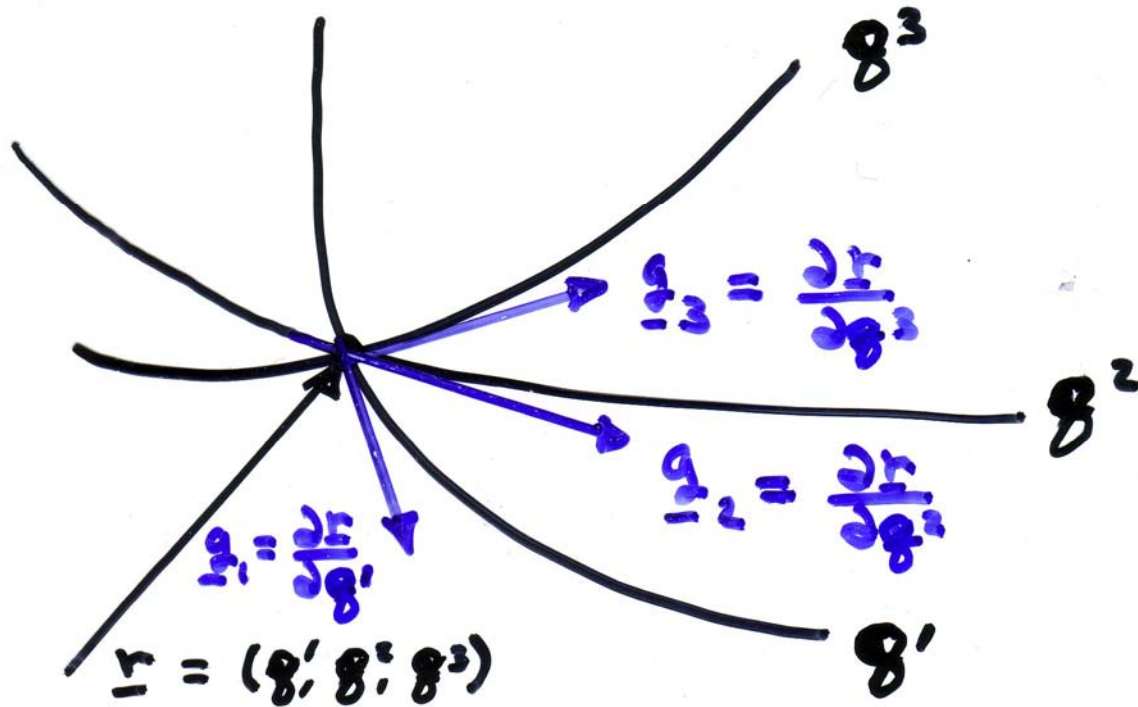


A.2 Curvilinear coordinates systems



Chain Rule

$$d\mathbf{r} = \underbrace{\frac{\partial \mathbf{r}}{\partial g^1}}_{\mathbf{g}_1} dg^1 + \underbrace{\frac{\partial \mathbf{r}}{\partial g^2}}_{\mathbf{g}_2} dg^2 + \underbrace{\frac{\partial \mathbf{r}}{\partial g^3}}_{\mathbf{g}_3} dg^3$$

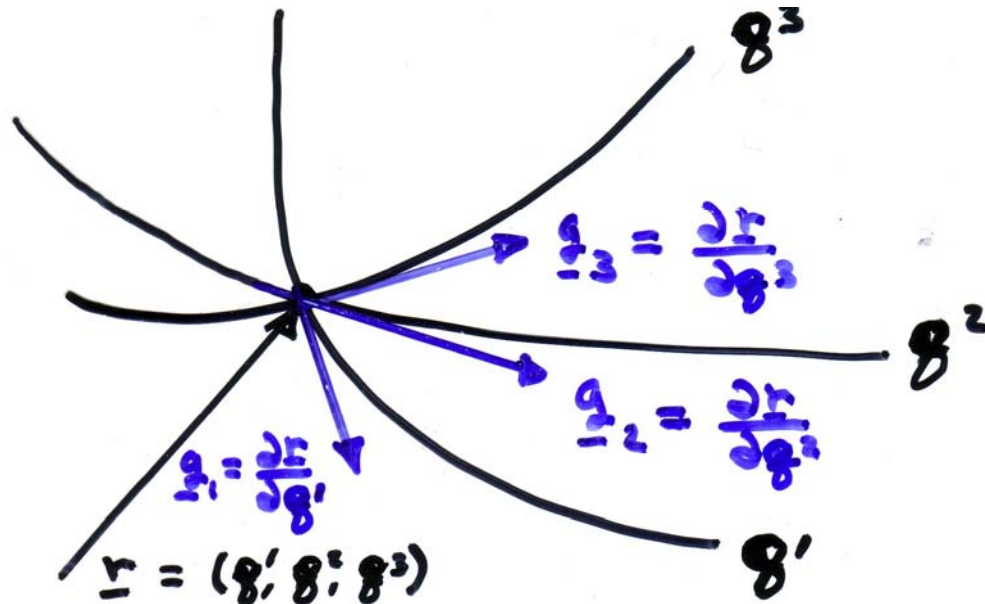
Base Vectors

$$\underline{g}_i = \frac{\partial \underline{r}}{\partial g^i} = \frac{\partial x^p}{\partial g^i} \underline{s}_p \quad (A2-4)$$

et

$$d\underline{r} = \underline{g}_1 dg^1 + \underline{g}_2 dg^2 + \underline{g}_3 dg^3 = \underline{g}_i dg^i$$

(Convention de \sum_i utilisée)



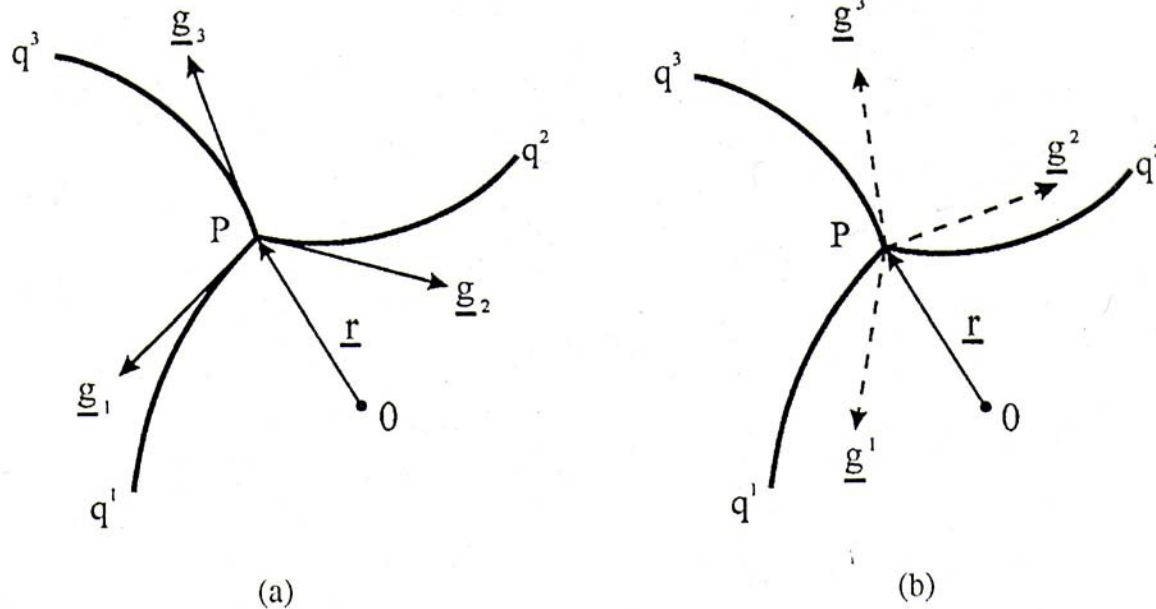


Figure A.2-2 Generalized coordinate system and base vectors: (a) covariant (tangential) base vectors; (b) contravariant (normal) base vectors

$$\underline{g}^n = \underline{\nabla} q^n = \text{grad } q^n = \frac{\partial q^n}{\partial x^l} \delta_l, \quad (\text{A.2-5a})$$

where the operator $\underline{\nabla}$ is defined by

$$\underline{\nabla} \equiv \underline{g}^n \frac{\partial}{\partial q^n}. \quad (\text{A.2-5b})$$

Note that neither \underline{g}_m nor \underline{g}^n are necessarily unit vectors, but from orthogonality it follows that

$$\underline{g}_m \cdot \underline{g}^n = \delta_m^n \quad (\text{A.2-6a})$$

Example A.2-1 Base Vectors in Spherical Coordinates

Soit (x^1, x^2, x^3) coordonnées Cartésiennes
et (r, θ, ϕ) coordonnées sphériques

From basic trigonometry

$$x^1 = r \sin \theta \cos \phi, \quad x^2 = r \sin \theta \sin \phi, \quad x^3 = r \cos \theta$$

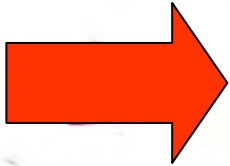
$$\begin{aligned} \Rightarrow \underline{g}_r &= \frac{\partial x^1}{\partial r} \underline{e}_1 + \frac{\partial x^2}{\partial r} \underline{e}_2 + \frac{\partial x^3}{\partial r} \underline{e}_3 \\ &= \sin \theta \cos \phi \underline{e}_1 + \sin \theta \sin \phi \underline{e}_2 + \cos \theta \underline{e}_3 \\ \underline{g}_\theta &= \frac{\partial x^1}{\partial \theta} \underline{e}_1 + \frac{\partial x^2}{\partial \theta} \underline{e}_2 + \frac{\partial x^3}{\partial \theta} \underline{e}_3 \\ &= r \cos \theta \cos \phi \underline{e}_1 + r \cos \theta \sin \phi \underline{e}_2 - r \sin \theta \underline{e}_3 \\ \underline{g}_\phi &= \frac{\partial x^1}{\partial \phi} \underline{e}_1 + \frac{\partial x^2}{\partial \phi} \underline{e}_2 + \frac{\partial x^3}{\partial \phi} \underline{e}_3 \\ &= -r \sin \theta \sin \phi \underline{e}_1 + r \sin \theta \cos \phi \underline{e}_2 \end{aligned}$$

From B.S.L. A.6

$$\underline{s}_r = \sin\theta \cos\phi \underline{s}_1 + \sin\theta \sin\phi \underline{s}_2 + \cos\theta \underline{s}_3$$

$$\underline{s}_\theta = \cos\theta \cos\phi \underline{s}_1 + \cos\theta \sin\phi \underline{s}_2 - \sin\theta \underline{s}_3$$

$$\underline{s}_\phi = -\sin\phi \underline{s}_1 + \cos\phi \underline{s}_2$$



$$\underline{g}_r = \underline{s}_r$$

$$\underline{g}_\theta = r \underline{s}_\theta$$

$$\underline{g}_\phi = r \sin\theta \underline{s}_\phi$$

Not necessarily
of unit value and
dimensionless!

To obtain the reciprocal (contravariant) base vectors

$$\mathbf{g}^i \cdot \mathbf{g}_j = \delta^i_j$$



$$\begin{aligned} \mathbf{g}_1 &= \mathbf{e}_1 \\ \mathbf{g}_2 &= \frac{\mathbf{e}_2}{\sqrt{2}} \\ \mathbf{g}_3 &= \frac{-\mathbf{e}_3}{\sqrt{2}} \end{aligned}$$



A.2-2 Transformation rule for Vectors

$$g^m = g^m(\bar{g}^n)$$

or $\bar{g}^n = \bar{g}^n(g^m)$

or $\bar{g}_m = \frac{\partial x^i}{\partial \bar{x}^m} \quad \text{or} \quad \bar{g}^n = \frac{\partial \bar{x}^n}{\partial x^i} \delta_{i2} = \nabla \bar{g}^n$

For any vector

$$\underline{v} = v^m \underline{g}_m = \bar{v}^n \bar{g}_n$$

$$\left(\frac{\partial x^i}{\partial \bar{x}^m} \right) \left(\frac{\partial \bar{x}^n}{\partial x^i} \right) \leftarrow \text{Chain Rule}$$

$$\underline{v} = v^m \bar{g}_n \left(\frac{\partial \bar{x}^n}{\partial x^m} \right) = \bar{v}^n \bar{g}_n$$

⇒

$$\bar{v}^n = v^m \left(\frac{\partial \bar{x}^n}{\partial x^m} \right)$$

Transformation of a contravariant component

N.B. For the coordinates:

$$\bar{g}^i = \bar{g}^i(g^1, g^2, g^3)$$

Use Chain Rule:

$$d\bar{g}^i = \underbrace{\frac{\partial \bar{g}^i}{\partial g^j}}_{\text{contravariant}} dg^j$$

Contravariant transformation

② Pour les vecteurs de base :

$$\begin{aligned}\underline{g}_m &= \frac{\partial \underline{g}}{\partial \bar{g}^n} \frac{\partial \bar{g}^n}{\partial g^m} \\ &= \bar{g}_m \frac{\partial \bar{g}^n}{\partial g^m}\end{aligned}$$

} transformation
d'une "composante"
covariante

Covariante covariante :

$$\underline{y} = v_m \underline{g}^m = \bar{v}_n \bar{g}^n$$

$$\nabla \underline{g}^m = \frac{\partial \underline{g}^m}{\partial \bar{g}^n} \nabla \bar{g}^n$$

⇒

$$v_m \frac{\partial \underline{g}^m}{\partial \bar{g}^n} \bar{g}^n = \bar{v}_n \bar{g}^n$$

et

$$\bar{v}_n = v_m \frac{\partial \underline{g}^m}{\partial \bar{g}^n}$$

Transformation
d'une composante
covariante

N.B. Transformation of reciprocal base vector

$$\underline{g}^m = \nabla \underline{g}^m = \frac{\partial \underline{g}^m}{\partial \bar{g}^n} \nabla \bar{g}^n = \frac{\partial \underline{g}^m}{\partial \bar{g}^n} \bar{g}^n$$



$$\underline{g}^m = \frac{\partial \underline{g}^m}{\partial \bar{g}^n} \bar{g}^n$$

se transforma como
una componente
contravariante

Transformation of second order tensor

$$\tau^{ij} = \left(\frac{\partial \bar{g}^i}{\partial g^k} \right) \left(\frac{\partial \bar{g}^j}{\partial g^l} \right) \bar{\tau}^{kl}$$

Componentes
contravariantes

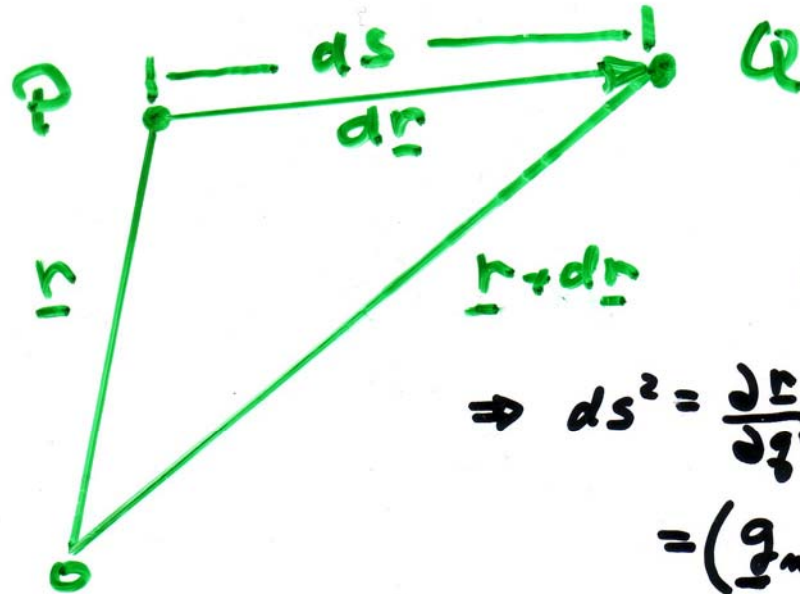
$$\tau_{ij} = \left(\frac{\partial \bar{g}^k}{\partial g^i} \right) \left(\frac{\partial \bar{g}^l}{\partial g^j} \right) \bar{\tau}_{kl}$$

Componentes
covariantes

$$\tau^i_j = \left(\frac{\partial \bar{g}^i}{\partial g^k} \right) \left(\frac{\partial \bar{g}^l}{\partial x^j} \right) \bar{\tau}^k_l$$

Componentes
mixtas.

A.2-4 Metric Tensor



$$ds^2 = dr \cdot dr$$

$$\Rightarrow ds^2 = \frac{\partial r}{\partial g^m} dg^m \cdot \frac{\partial r}{\partial g^n} dg^n$$

$$= \underbrace{(g_m \cdot g_n)}_{g_{mn}} dg^m dg^n$$



$$ds^2 = \underbrace{g_{mn}} dg^m dg^n$$

Covariant components of metric tensor

Properties:

N.B. ① en coord. orthogonales $g_{mn} = \delta_{mn}$

ou général $g_{mn} \neq \delta_{mn}$

② $g_{mn} = g_{nm}$ (symétrique)

③ $\underline{g}^r \cdot \underline{g}^s = g^{rs}$ *composantes contravariantes du tenseur métrique.*

④ $g_{mj} g^{jn} = (g_m \cdot \underline{g}_j) (\underline{g}^j \cdot \underline{g}^n)$
 $= (g_m \cdot \underline{g}^j) (\underline{g}_j \cdot \underline{g}^n)$
 $= \delta_m^j \delta_j^n = \delta_m^n$

A.2-5 Physical Components

On a: $\underline{v} = v_m \underline{g}^m = v^n \underline{g}_n$

$\cdot \underline{g}_r$ $v_m \underline{g}^m \cdot \underline{g}_r = v^n \underline{g}_n \cdot \underline{g}_r$

$v_m \delta_r^m = v^n g_{nr}$

\Rightarrow

$$v_r = g_{nr} v^n$$

\underline{e}_n

Transformation
de comp. cartés.
à comp. covar.

Physical component of the base vectors

$$g_{(n)} = \frac{g_n}{|g_n|} = \frac{g_n}{\sqrt{g_{nn}}}$$

No
summation!

Physical components of a vector

$$v_{(n)} = \sqrt{g_{nn}} v^n = \sqrt{g_{nn}} g^{nm} v_m = \sqrt{g^{nn}} v_n$$

Physical components of a second order tensor

$$T_{(mr)} = \sqrt{g_{mm}} \sqrt{g_{rr}} T^{mr} = \sqrt{g^{mm}} \sqrt{g^{rr}} T_{mr} = \sqrt{g^{mm}} \sqrt{g_{rr}} T_m^r$$

Etc...

Examples A.2-4 & 5: metric tensor and physical components in spherical coordinates

$$g_{ij} = \underline{g}_i \cdot \underline{g}_j \qquad g^{ij} = \underline{g}^i \cdot \underline{g}^j$$

$$g_{rr} = \underline{g}_r \cdot \underline{g}_r = 1 = g^{rr}$$

$$g_{\theta\theta} = r \underline{g}_\theta \cdot r \underline{g}_\theta = r^2 = (1/g^{00})^2 \leftarrow$$

$$g_{\phi\phi} = r \sin\theta \underline{g}_\phi \cdot r \sin\theta \underline{g}_\phi = r^2 \sin^2\theta = (1/g^{11})^2 \leftarrow$$

$$g_{ij} = 0 \text{ for } i \neq j \qquad 0 = g^{ij} \quad j \neq i$$

$$g_{ij} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & r^2 & 0 \\ 0 & 0 & r^2 \sin^2\theta \end{pmatrix}, \quad g^{ij} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & r^{-2} & 0 \\ 0 & 0 & r^{-2} \sin^{-2}\theta \end{pmatrix}$$

Vector

$$\begin{aligned}\underline{V} &= V_r \underline{\hat{e}}_r + V_\theta \underline{\hat{e}}_\theta + V_\phi \underline{\hat{e}}_\phi \\ &= V_1 \underline{\hat{e}}^1 + V_2 \underline{\hat{e}}^2 + V_3 \underline{\hat{e}}^3 \\ &= V^1 \underline{\hat{e}}_1 + V^2 \underline{\hat{e}}_2 + V^3 \underline{\hat{e}}_3\end{aligned}$$

$$\Rightarrow V^{(1)} = \sqrt{g_{11}} V^1 = \sqrt{g^{11}} V_1 = V_r = V_1 = V^1$$

$$V^{(2)} = \sqrt{g_{22}} V^2 = \sqrt{g^{22}} V_2 = V_\theta$$

$$V^2 = V_\theta / r \quad \text{at} \quad V_2 = r V_\theta$$

$$V^{(3)} = \sqrt{g_{33}} V^3 = \sqrt{g^{33}} V_3 = V_\phi$$

$$V^3 = \frac{V_\phi}{r \sin \theta} \quad \text{at} \quad V_3 = r \sin \theta V_\phi$$

A.3 Covariant differentiation

Operator DEL, ∇ :

$$\nabla \equiv \underline{g}^i \frac{\partial}{\partial g^i}$$

Gradient of a scalar

$$\nabla s = \underline{g}^i \frac{\partial s}{\partial g^i} = \underline{g}^i s_i$$

Gradient of a vector

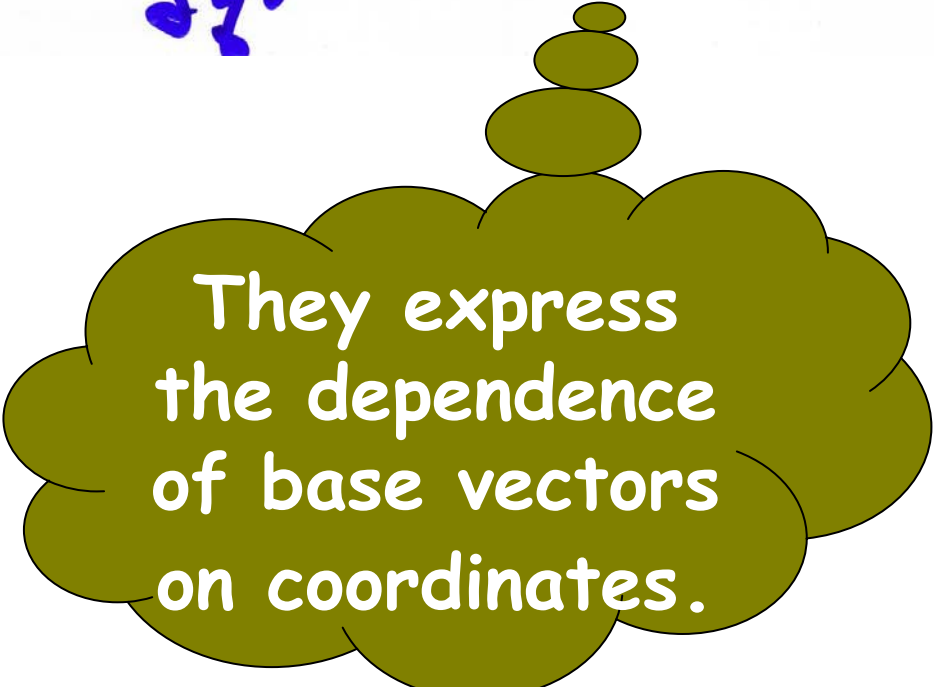
$$\begin{aligned} \nabla \underline{v} &= \underline{g}^i \frac{\partial}{\partial g^i} (\underline{g}^j v_j) \\ &= \underline{g}^i \underline{g}^j \frac{\partial v_j}{\partial g^i} + \underline{g}^i v_j \underbrace{\frac{\partial}{\partial g^i} \underline{g}^j}_{-\Gamma_{ki}^j \underline{g}^k} \end{aligned}$$

Christoffel symbols

Christoffel symbols

$$\frac{\partial \underline{g}_i}{\partial g_j} = \Gamma_{ij}^k \underline{g}_k = \{ \begin{smallmatrix} k \\ ij \end{smallmatrix} \} \underline{g}_k$$

$$\frac{\partial \underline{g}^i}{\partial g_j} = -\Gamma_{kj}^i \underline{g}^k$$



They express
the dependence
of base vectors
on coordinates.

Properties:

$$\Gamma_{mj}^l = \frac{1}{2} g^{lk} \left[\frac{\partial g_{ik}}{\partial g^m} + \frac{\partial g_{mk}}{\partial g^j} - \frac{\partial g_{mj}}{\partial g^k} \right] \quad (A.3-10)$$

$$\Gamma_{mj}^l = \Gamma_{jm}^l \quad \text{symétrique}$$

$$\begin{aligned} \nabla_{\underline{v}} &= \underline{g}^i \underline{g}^j \frac{\partial}{\partial g^i} v_j - \underline{g}^i \underline{g}^k v_j \Gamma_{ki}^j \\ &= \underline{g}^i \underline{g}^j \left(\frac{\partial v_j}{\partial g^i} - v_k \Gamma_{ji}^k \right) \end{aligned}$$

$$\rightarrow \boxed{\nabla_{\underline{v}} = \underline{g}^i \underline{g}^j v_{j,i}}$$

Also:

$$\underline{\nabla} \underline{v} = \underline{g}^i \underline{g}_j v_{,i}^j$$

$$\underline{\nabla} \underline{\underline{\tau}} = \underline{g}^i \underline{g}^j \underline{g}^k \tau_{jk, i}$$

$$\tau_{jk, i} = \frac{\partial}{\partial q^i} \tau_{jk} - \Gamma_{ji}^l \tau_{lk} - \Gamma_{ki}^l \tau_{jl}$$

Divergence of a vector:

$$\underline{\nabla} \cdot \underline{v} = \underline{g}^i \frac{\partial}{\partial q^i} \cdot v_j \underline{g}^j = g^{ij} \frac{\partial v_j}{\partial q^i} - g^{ik} v_j \Gamma_{ki}^j$$

$$= g^{ij} \left(\frac{\partial v_j}{\partial q^i} - v_k \Gamma_{ji}^k \right) = g^{ij} v_{j,i}$$

$$\Rightarrow \underline{\nabla} \cdot \underline{v} = v^i_{,i} \text{ a scalar}$$